

Lab 8

Your assignment this week will be another very non-trivial generalization of your code. I will try to explain what you are to do as clearly as possible, and give you some hints, as appropriate.

Because this will be a fair bit of work (probably even more than Lab 7), the due date for this lab will be Wednesday, 01-Apr-2009, at 5PM.

Your assignment is as follows:

- To make it easier for me to run your codes, you are to take all of your input from my input deck, just like last week. By “input”, I mean all of the data for the PDE (i.e., $\mathbf{k}(\mathbf{x})$, $r(\mathbf{x})$, BCs), as well as the FE parameters (number of elements, element type, etc.). I have posted it on bspace. It is fairly well documented, but please ask if anything is unclear.

The deck is just a matlab function that you can (i.e. *will*) call, and get back (hopefully) all of the data needed to run your code. Please, do not have your code take any command line input; get everything from the input deck.

- Solve a heat transfer problem in 2D. Recall from Homework 5 the exact weak form, with the modification that \mathbf{k} is both anisotropic, and a function of \mathbf{x} :

$$\begin{aligned}
 0 &= \int_{\Omega} [\operatorname{div}(\mathbf{k}(\mathbf{x}) \operatorname{grad}(T)w) - \mathbf{k}(\mathbf{x}) \operatorname{grad}(T) \cdot \operatorname{grad}(w) + wr] dv, \\
 &\Rightarrow \\
 \int_{\Omega} -\mathbf{k}(\mathbf{x}) \operatorname{grad}(T) \cdot \operatorname{grad}(w) dv &= - \int_{\Omega} wr dv - \int_{\Omega} \operatorname{div}(\mathbf{k}(\mathbf{x}) \operatorname{grad}(T)w) dv, \\
 \int_{\Omega} \operatorname{grad}(w) \cdot \mathbf{k}(\mathbf{x}) \operatorname{grad}(T) dv &= \int_{\Omega} wr dv - \int_{\partial\Omega} w\mathbf{q} \cdot \mathbf{n} da.
 \end{aligned}$$

Also, recall its approximation:

$$\begin{aligned}
 \int_{\Omega} (\mathbf{B}\hat{\mathbf{w}}) \cdot \mathbf{k}(\mathbf{x}) (\mathbf{B}\hat{\mathbf{T}}) dv &= \int_{\Omega} (\mathbf{N}\hat{\mathbf{w}}) r dv - \int_{\partial\Omega} \bar{q} (\mathbf{N}\hat{\mathbf{w}}) da, \\
 \hat{\mathbf{w}}^T \left(\left[\int_{\Omega} \mathbf{B}^T \mathbf{k}(\mathbf{x}) \mathbf{B} dv \right] \hat{\mathbf{T}} \right) &= \hat{\mathbf{w}}^T \left(\int_{\Omega} \mathbf{N}^T r dv - \int_{\partial\Omega} \bar{q} \mathbf{N}^T da \right),
 \end{aligned}$$

where

$$\mathbf{k} = \begin{bmatrix} k_{1,1}(\mathbf{x}) & k_{1,2}(\mathbf{x}) \\ k_{2,1}(\mathbf{x}) & k_{2,2}(\mathbf{x}) \end{bmatrix}.$$

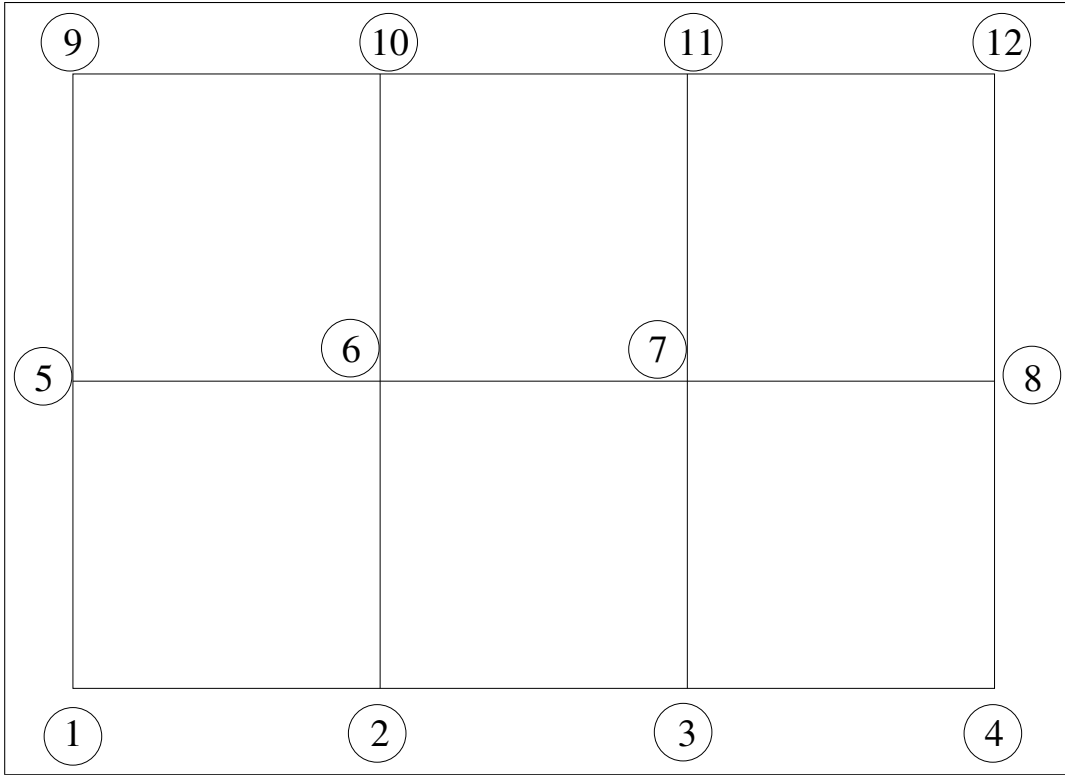
Thus,

$$\left[\int_{\Omega} \mathbf{B}^T \mathbf{k}(\mathbf{x}) \mathbf{B} dv \right] \hat{\mathbf{T}} = \int_{\Omega} \mathbf{N}^T r dv - \int_{\partial\Omega} \bar{q} \mathbf{N}^T da,$$

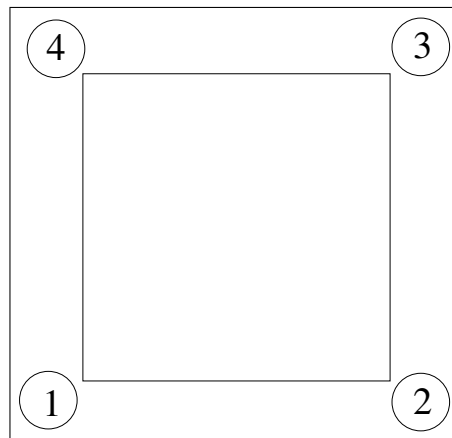
or

$$\mathbf{K} \hat{\mathbf{T}} = \mathbf{F}.$$

- You can assume that the domain is always a rectangle in \mathbb{R}^2 , such that $x_1 \in \{x : 0 \leq x \leq a\}$, and $x_2 \in \{x : 0 \leq x \leq b\}$.
- You can assume that the global node numbering will always be monotonically increasing, *first* in the x_1 direction, and then in the x_2 direction. An illustration follows, for the case of $nel_1 = 3$ and $nel_2 = 2$.



Note that the *local* node numbering for 4 node quads is typically given as follows:



- Your elements only need to be linear quadrilaterals, formulated in parametric space.
- In the 1D case, many of you used the chain rule to claim that

$$\frac{\partial N_A}{\partial x} \frac{\partial N_B}{\partial x} \frac{\partial x}{\partial \xi} = \frac{\partial N_A}{\partial \xi} \frac{\partial N_B}{\partial \xi} \frac{\partial \xi}{\partial x},$$

which is ok for the 1D case, but won't work for 2D or 3D. So, we have to calculate $\frac{\partial N}{\partial \mathbf{x}}$ properly. First, note that we represent the gradient of a scalar T in 2D as follows:

$$\text{grad}(T) \approx \mathbf{B}\hat{\mathbf{T}},$$

where

$$\mathbf{B} = \begin{bmatrix} N_{1,1} & N_{2,1} & \dots & N_{nen,1} \\ N_{1,2} & N_{2,2} & \dots & N_{nen,2} \end{bmatrix}.$$

So, if we can calculate $\frac{\partial N_A}{\partial x}$ and $\frac{\partial N_A}{\partial y}$ for all A , then we are good. So, how to do this? We don't have the shape functions written in physical space anymore, so we can't take those derivatives directly, but we also know that there is some relation between points on the master element, and points in the physical element, and we exploit that fact via the chain rule:

$$\begin{aligned} \begin{bmatrix} \frac{\partial N_A}{\partial \xi} \\ \frac{\partial N_A}{\partial \eta} \end{bmatrix} &= \begin{bmatrix} \frac{\partial N_A}{\partial x} \frac{\partial x}{\partial \xi} + \frac{\partial N_A}{\partial y} \frac{\partial y}{\partial \xi} \\ \frac{\partial N_A}{\partial x} \frac{\partial x}{\partial \eta} + \frac{\partial N_A}{\partial y} \frac{\partial y}{\partial \eta} \end{bmatrix}, \\ &= \begin{bmatrix} \frac{\partial x}{\partial \xi} & \frac{\partial y}{\partial \xi} \\ \frac{\partial x}{\partial \eta} & \frac{\partial y}{\partial \eta} \end{bmatrix} \begin{bmatrix} \frac{\partial N_A}{\partial x} \\ \frac{\partial N_A}{\partial y} \end{bmatrix}. \end{aligned}$$

So, the term on the left hand side is easy to calculate, since the shape functions are formulated directly in isoparametric coordinates. The matrix on the right hand side is also easy to calculate, for the same reason. The vector on the right hand side is what we want to solve for, so, solve the system of equations.

- Your code should perform its integration using Gaussian quadrature. Note that the matrix on the right hand side above is what we call the “Jacobian” matrix for the element, and its determinant, $J = \det([J])$, is the scale factor we use to map the integrals to the parent space. Thus, the quadrature (of the stiffness, for example) can be calculated as follows:

$$\begin{aligned}
\int_{\Omega^e} \mathbf{B}^T \mathbf{k}(\mathbf{x}) \mathbf{B} \, dv &= \int_{-1}^1 \int_{-1}^1 \mathbf{B}(\xi, \eta)^T \mathbf{k}(\mathbf{x}(\boldsymbol{\xi})) \mathbf{B}(\xi, \eta) J(\xi, \eta) \, d\xi \, d\eta, \\
&= \int_{-1}^1 \int_{-1}^1 g(\xi, \eta) J(\xi, \eta) \, d\xi \, d\eta, \\
&\approx \int_{-1}^1 \left(\sum_{i=1}^{nip} g(\xi_i, \eta) J(\xi_i, \eta) w_i \right) d\eta, \\
&\approx \sum_{j=1}^{nip} w_j \left(\sum_{i=1}^{nip} g(\xi_i, \eta_j) J(\xi_i, \eta_j) w_i \right), \\
&= \sum_{i=1}^{nip} \sum_{j=1}^{nip} g(\xi_i, \eta_j) J(\xi_i, \eta_j) w_i w_j.
\end{aligned}$$

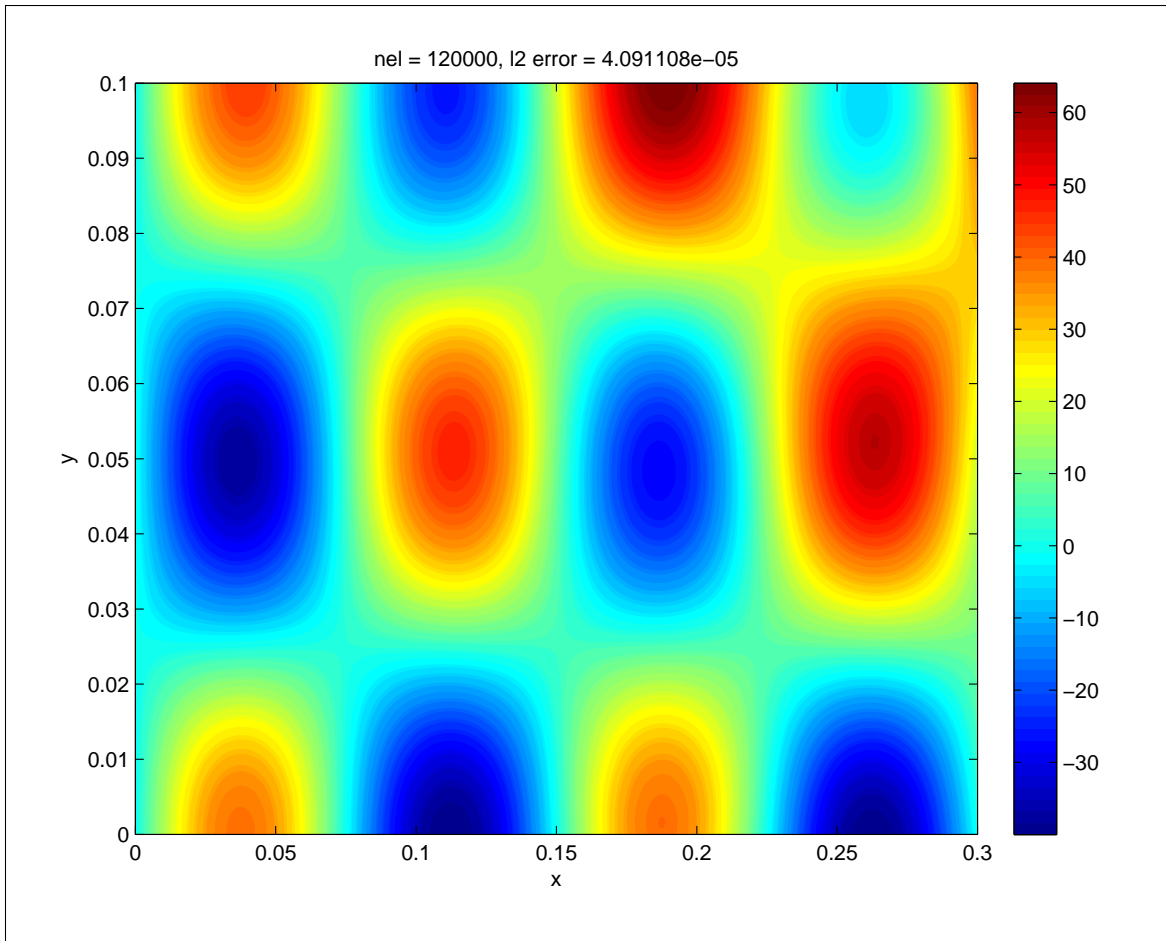
- You only need to implement Dirichlet boundary conditions. Implementation of Neumann boundary conditions would be a bit complicated, and I figure that, ultimately, you already know how to do integrations in 1D . . .

Per the input deck, you can assume that each Dirichlet BC covers a whole edge, and the edges are numbered, as follows:

edge number	description
1	$x = 0$
2	$x = a$
3	$y = 0$
4	$y = b$

You need to turn the following in to me:

- A convergence plot of your code versus the exact solution, in the l^2 norm, for the problem given in Homework 6. Plot ten different pairs of (nel_1, nel_2) , ensuring that your elements are always square, and such that your last run meets the same tolerance given in Homework 6. Plot your error versus $\frac{1}{h^e}$. Comment on how many linear elements your code took to meet the stated tolerance, versus COMSOL (from Homework 6).
- A plot of the output from your final run. It should look something like the following:



- A plot of the error $|u_{exact} - u_{FE}|$ as a function of x and y .

Please give these to me on paper (not via email).

- Email me your code, which I will run with an input deck that is different from the one I posted.